

I. Selection rules for IR transitions**A. IR transitions (Harmonic oscillator wave functions):**

Often it is possible to see that a particular integral $\int \psi_m s \psi_n ds = 0$ on the basis of symmetry.

- For example, for $m = 0$ and $n = 2$, sketch the wave functions, multiply ψ_0 first by s , next by ψ_2 and then plot the result. Use this plot to show that the integral over all values of s must be zero.
- Plot $\psi_m s \psi_n ds$ for $m = 3, n = 5$ over a reasonable range of s values, and then make a statement about total area under the curve based upon the symmetry of your plot.
- Sometimes the symmetry argument will not work. Plot $\psi_m s \psi_n ds$ for $m = 2, n = 5$ and then use numerical integration to show that the integral over a large range of s values is close to 0.
- There is an important relationship between successive Hermite polynomials which is worked out on p. 79 of Pauling and Wilson:

$$H_{n+1}(s) - 2sH_n(s) + 2nH_{n-1}(s) = 0.$$

This formula allows one to substitute for sH_n in the integral $\int \psi_m s \psi_n ds$ and then use the orthogonality of ψ_m and ψ_n to derive a general rule for state transitions in a harmonic oscillator. Try doing this to show that $\int \psi_m s \psi_n ds = 0$ unless $m = n+1$ or $m = n-1$. This result is known as a selection rule.

- Use the formula of problem 4 to show that in general,

$$s^k H_n = a_{n+k} H_{n+k} + a_{n+k-2} H_{n+k-2} + \dots + a_{n-k+2} H_{n-k+2} + a_{n-k} H_{n-k}$$

(if the subscript is < 0 , then the coefficient is 0)

- Try finding a "selection rule" for the integrals: $\int \psi_m s^2 \psi_n ds$ and $\int \psi_m s^3 \psi_n ds$ (Use the formula given in problem 5)
- Here is the relevance of the previous problem: Consider the perturbed HO, with $V(x)$ given by a Morse-like potential: $V(x) = \frac{1}{2} kx^2 + a_3 x^3 + \dots$ where the last term may be considered as a perturbing potential:

$$v(x) = a_3 x^3 + \dots = a' s^3 + \dots$$

This perturbing potential introduces modifications $\psi_m^{(1)}$ into the wavefunctions given by equation (39)

$$\text{from the perturbation theory section: } \psi_m^{(1)} = \sum_n \left[\frac{\langle \psi_n^{(0)*} | v | \psi_m^{(0)} \rangle}{E_m^{(0)} - E_n^{(0)}} \right] \psi_n^{(0)} = a' \sum_n \left[\frac{\langle \psi_n^{(0)*} | s^3 | \psi_m^{(0)} \rangle}{E_m^{(0)} - E_n^{(0)}} \right] \psi_n^{(0)} \text{ for } m \neq n$$

If a summation term is not 0 for any particular n , then this formula "mixes" $\psi_n^{(0)}$ into $\psi_m^{(0)}$.

- Show that the $\psi_m^{(1)}$ correction "mixes" states $m+3, m+1, m-1$, and $m-3$ into $\psi_m^{(0)}$
- Given this result, show that IR transitions with $n = m + 2$ or $m - 2$ are now "allowed". (If the a' term of the perturbation $a's^3$ is quite small, the transition probabilities are still very low).

B. Electronic transitions:

- Use symmetry arguments to show that $\int \psi_{1s} x \psi_{2s} d\tau = 0$ and that $\int \psi_{1s} x \psi_{2p_x} d\tau$ will not be zero. (An extension of this leads to the selection rule $\Delta l = \pm 1$ for electronic transitions.)